

# Magnetic field decay and unification of young and millisecond pulsar populations

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June 8, 2016

Physics of Pulsar Magnetospheres  
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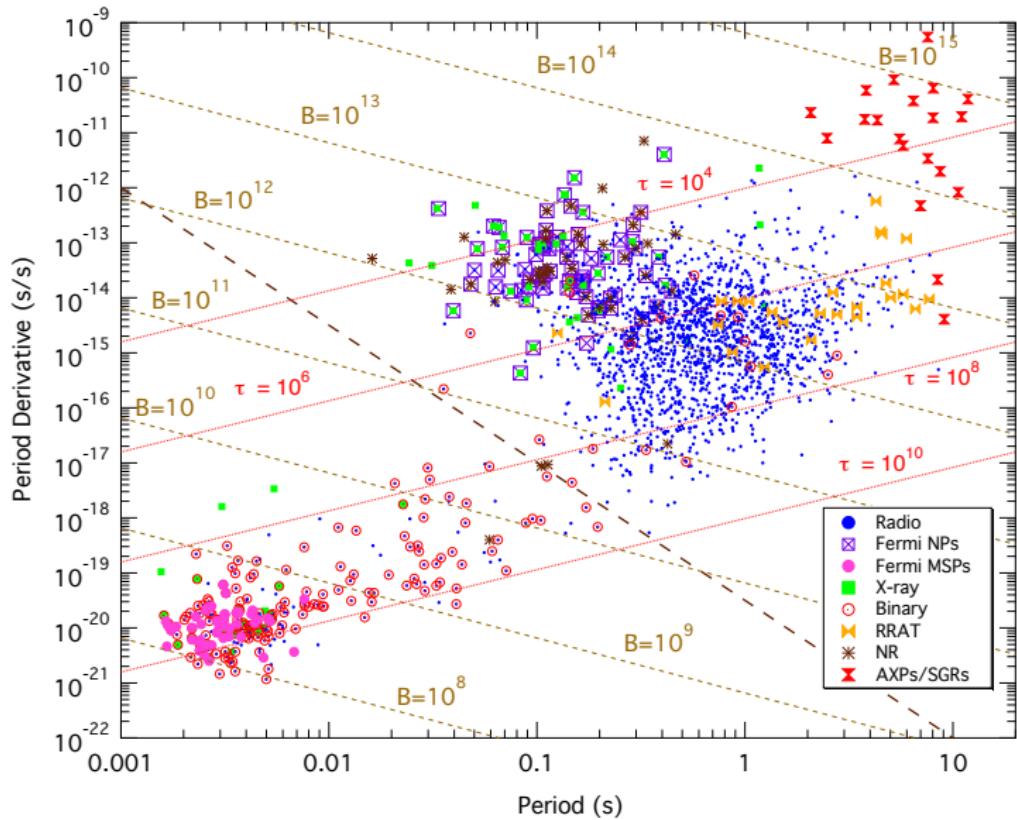
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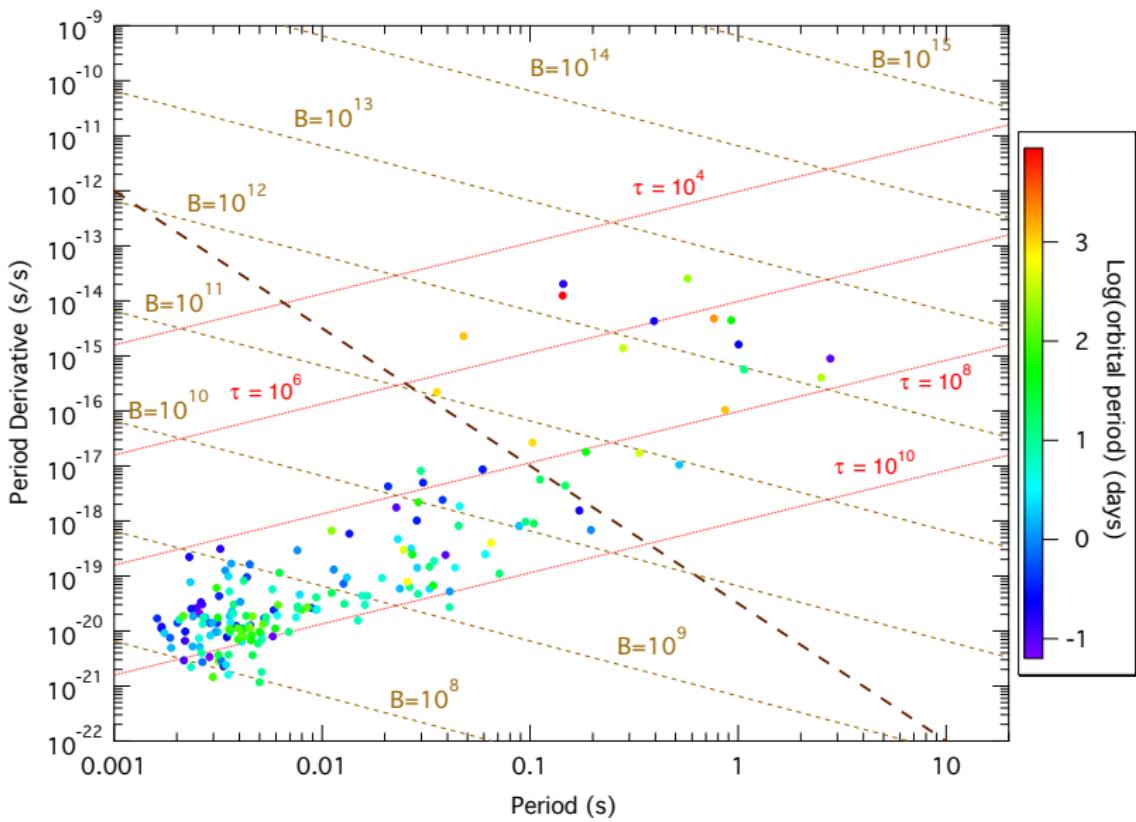
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- ① Pulsar zoo
- ② Population synthesis of normal pulsars
- ③ Population synthesis of millisecond pulsars
- ④ The Alicante connection: NPs to MSPs - field decay
- ⑤ Proposed new simulation
- ⑥ Conclusions

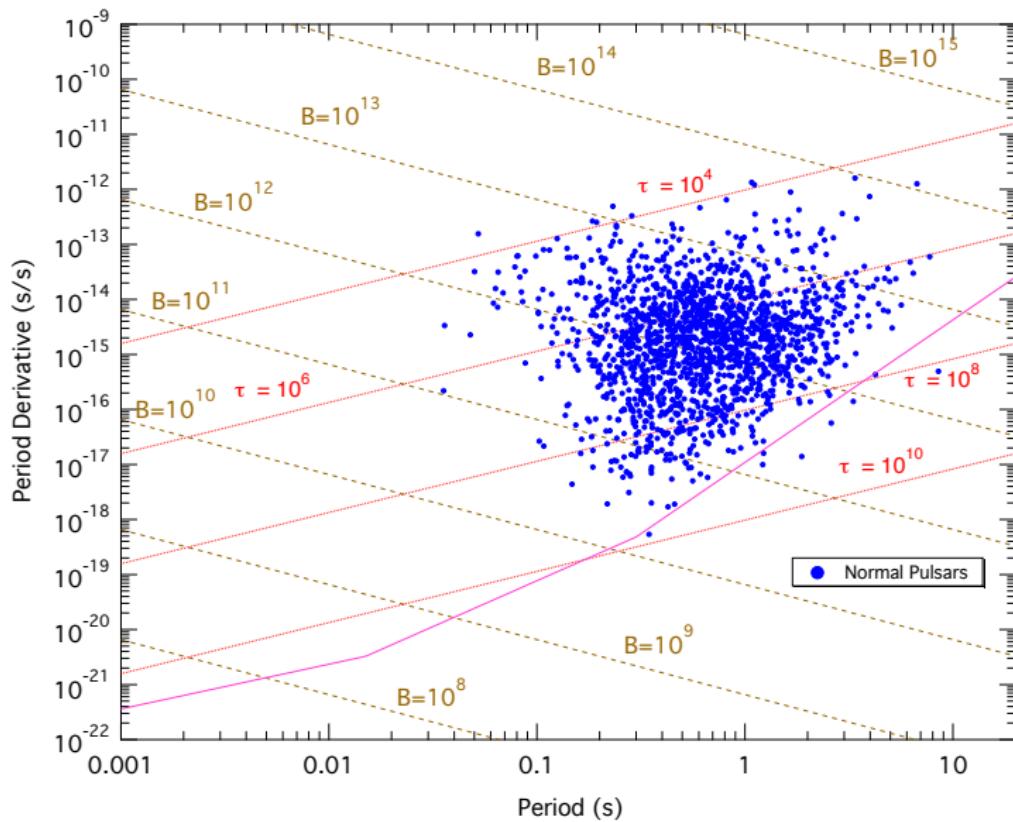
# Pulsar Zoo with over 2530 subjects



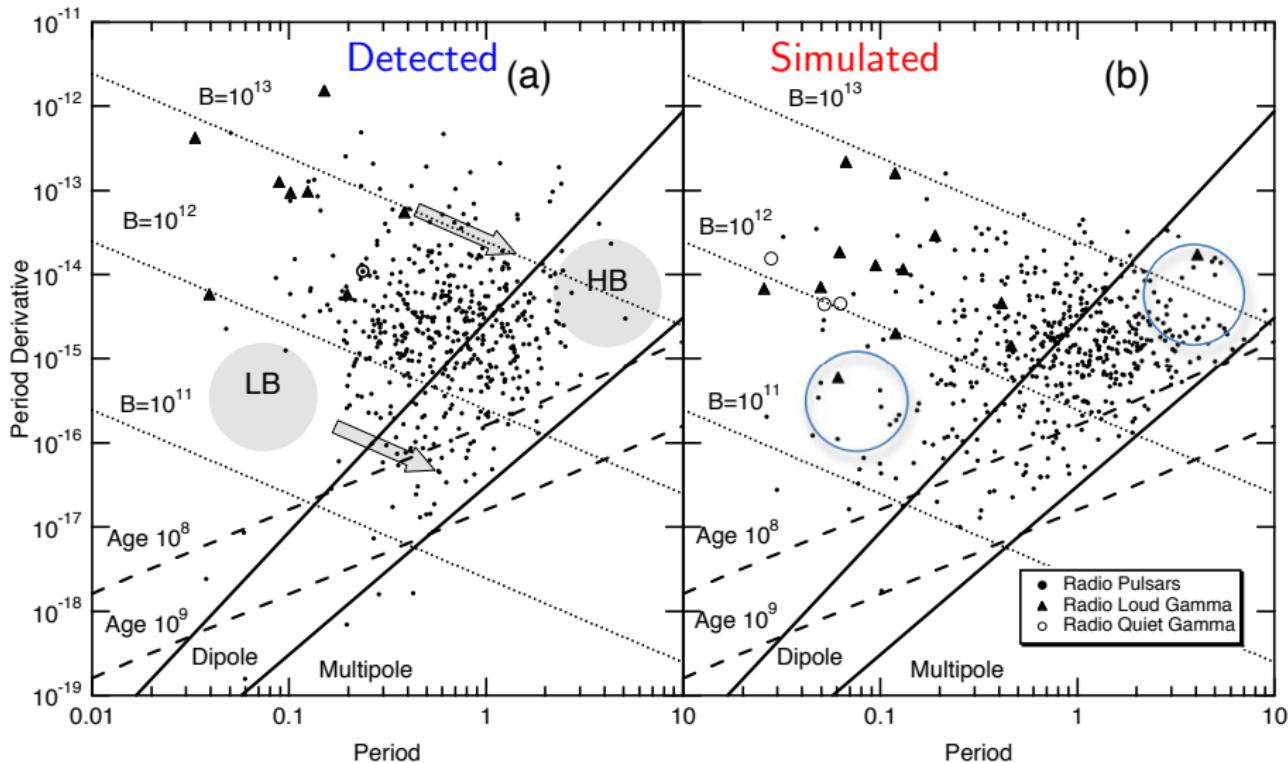
# Binary Pulsars



# Normal Pulsars

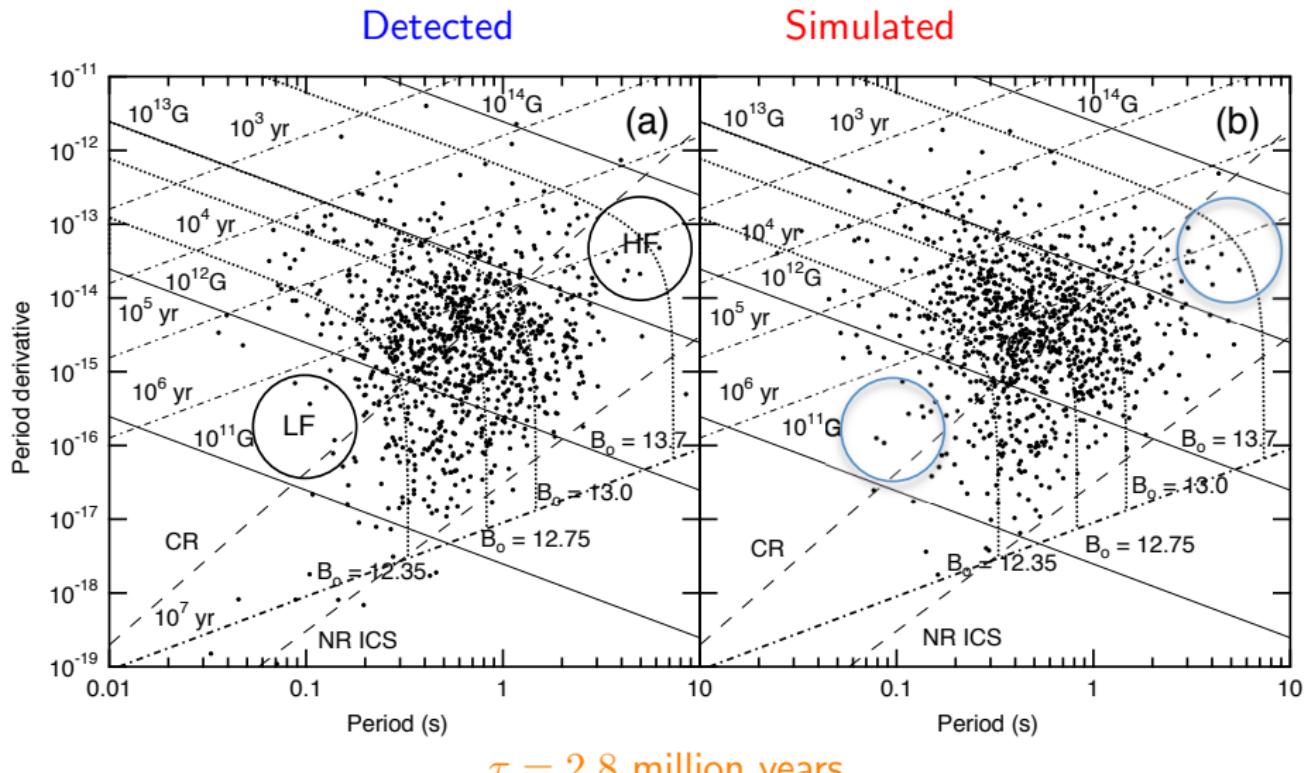


# Gonthier et al. (2002) — no field decay



Death lines - Zhang, Harding & Muslimov (2000)

# Gonthier et al. (2004) — Exponential B field decay



Death lines - Harding, Muslimov & Zhang (2002)

## Birth distributions — a more recent study

- Born in the spiral arms using the electron density model NE2001 — Cordes & Lazio (2003). Spiral arm system rotates with the speed of the density waves
- Trajectories are evolved in the Galactic potential — Paczyński (1990)
- For  $B_o$  — Log normal distribution

$$\langle \log B_o \rangle = 12.8$$

$$\sigma_{\log B_o} = 0.46$$

- For  $P_o$  — Gaussian distribution — need to explore other distributions

$$\langle P_o \rangle = 0.11 \text{ with } P_{o\min} \text{ of } 1.3 \text{ ms}$$

$$\sigma_{P_o} = 0.14$$

- Four free parameters to define the means and widths — searched in MCMC chains

## Spin-down and field decay

- Spitkovsky (2006)
- Contopoulos, Kalapotharakos & Kazanas (2014)
- Tchekhovskoy, Philippov & Spitkovsky (2016)

$$L = (1 + \sin^2 \chi) L_o$$

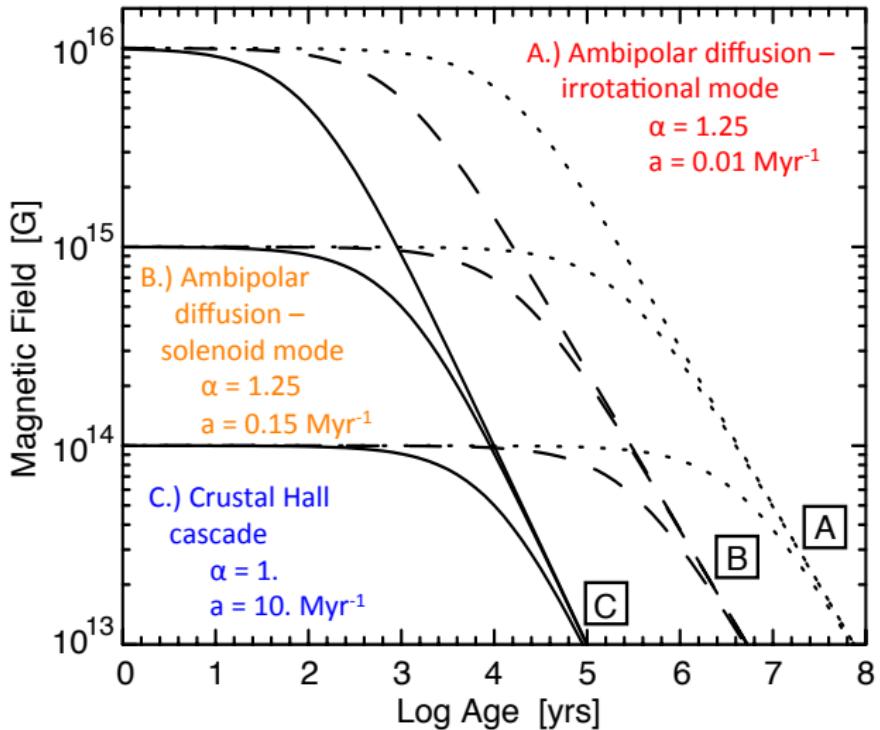
- Magnetic field decay model — Colpi, Geppert & Page (2000)

$$\frac{d B_{13}}{d t} = -a B_{13}^{1+\alpha} \quad (\text{Eq. 2})$$

$$a = \frac{1}{\tau_{\text{Myr}}}$$

$$B_{13}(t) = \frac{B_{13_o}}{\left[1 + a \alpha B_{13_o}^\alpha t_6\right]^{1/\alpha}} \quad (\text{Eq. 3})$$

# Colpi, Geppert, & Page 2000 - Figure 1



Colpi, Geppert, & Page 2000, ApJ, 529, L29 – Figure 1

# Magnetic field decay — Viganò, Pons, & Miralles 2012

Viganò et al. (2013) —  
Figure 10

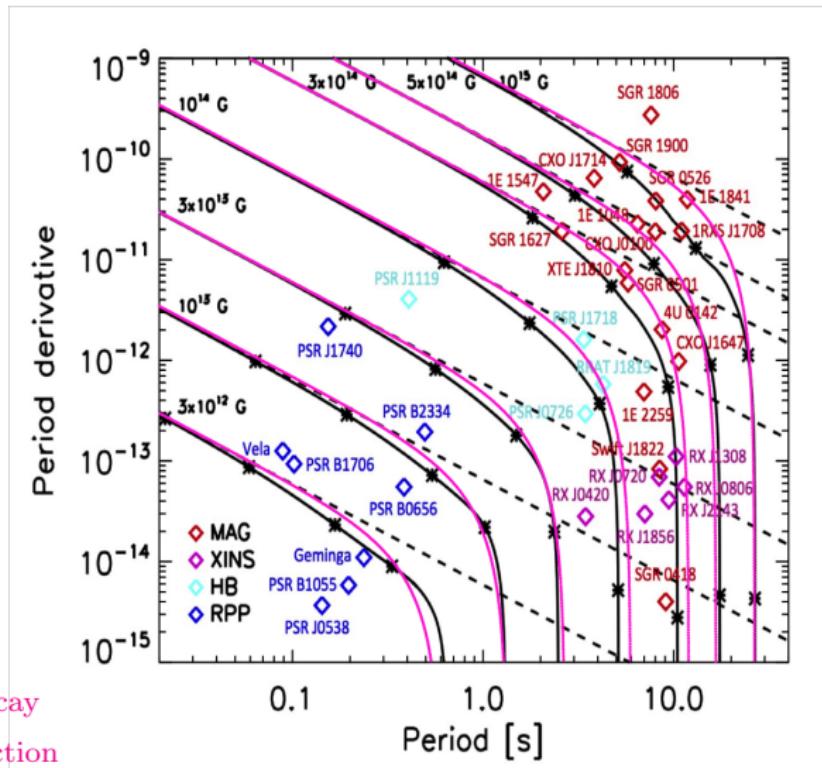
$$a = 1 \text{ Myr}^{-1}$$

$$\tau_{\text{Myr}} = 1 \text{ Myr}$$

$$\alpha = 0.7$$

$$\frac{d B_{13}}{dt} = -a B_{13}^{1+\alpha}$$

$$\alpha = \begin{cases} 0 \rightarrow \text{Ohmic decay} \\ 1 \rightarrow \text{Hall induction} \\ 2 \rightarrow \text{Ambipolar diffusion} \end{cases}$$



private communication, José Pons

# Radio and $\gamma$ -ray beam geometry and emission

- Harding, Grenier & Gonthier (2007) and
- Pierbattista, Grenier, Harding, & Gonthier (2012)
- Core and conal beams
  - Conal — altitude dependent — Kijak & Gil (2003)
  - Empirical radio luminosity —  $P$  and  $\dot{P}$  dependence Arzoumanian, Chernoff, Cordes (2002)
$$\mathcal{L}_\nu = \mathcal{L}_0 P^{\alpha_\nu} \dot{P}^{\beta_\nu} \text{ mJy} \cdot \text{kpc}^2 \cdot \text{MHz}$$
- Exponents  $\alpha_\nu$  and  $\beta_\nu$  are free parameters searched by MCMC

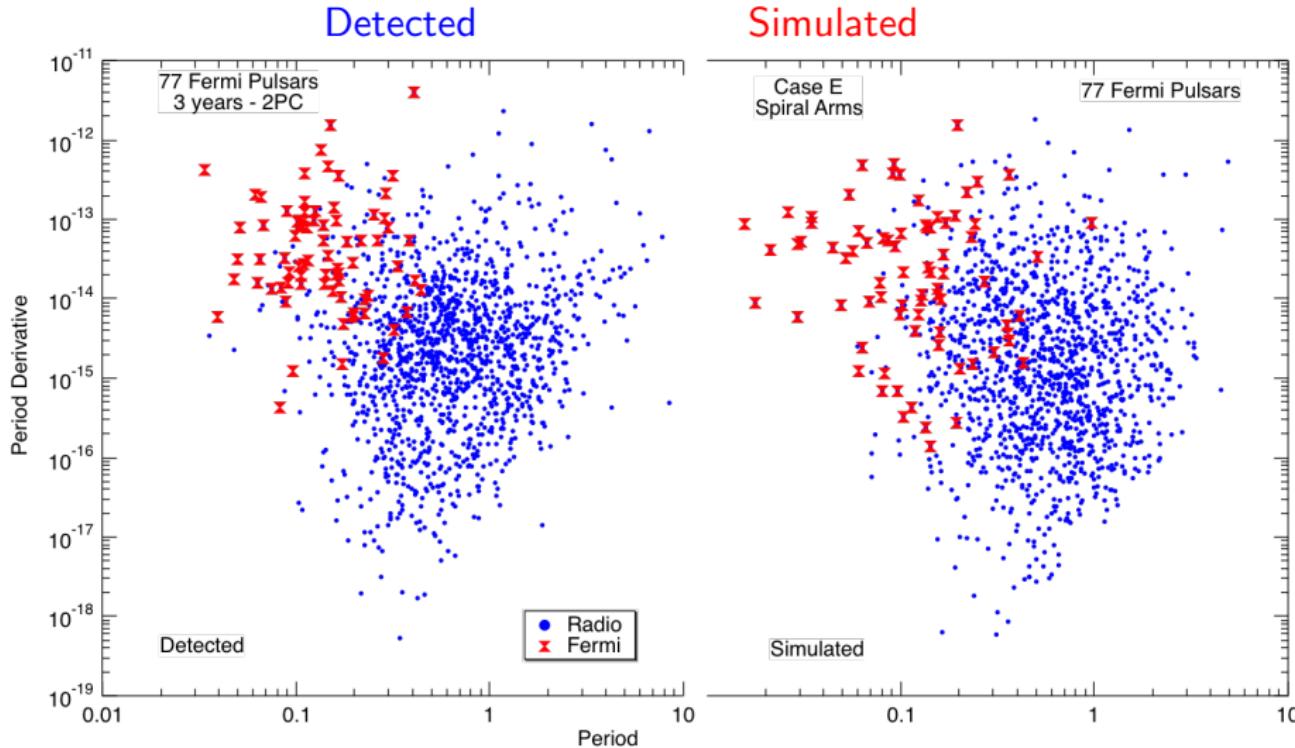
$$\alpha_\nu = -0.94, \quad \beta_\nu = 0.41$$

- Threshold characteristics of a select group of ten radio surveys
- $\gamma$ -ray sky maps — Extended Slot Gap emission — Muslimov & Harding (2004)
  - Empirical  $\gamma$ -ray luminosity

$$\mathcal{L}_\gamma = f_\gamma P^{\alpha_\gamma} \dot{P}^{\beta_\gamma} \text{ eV/s}$$

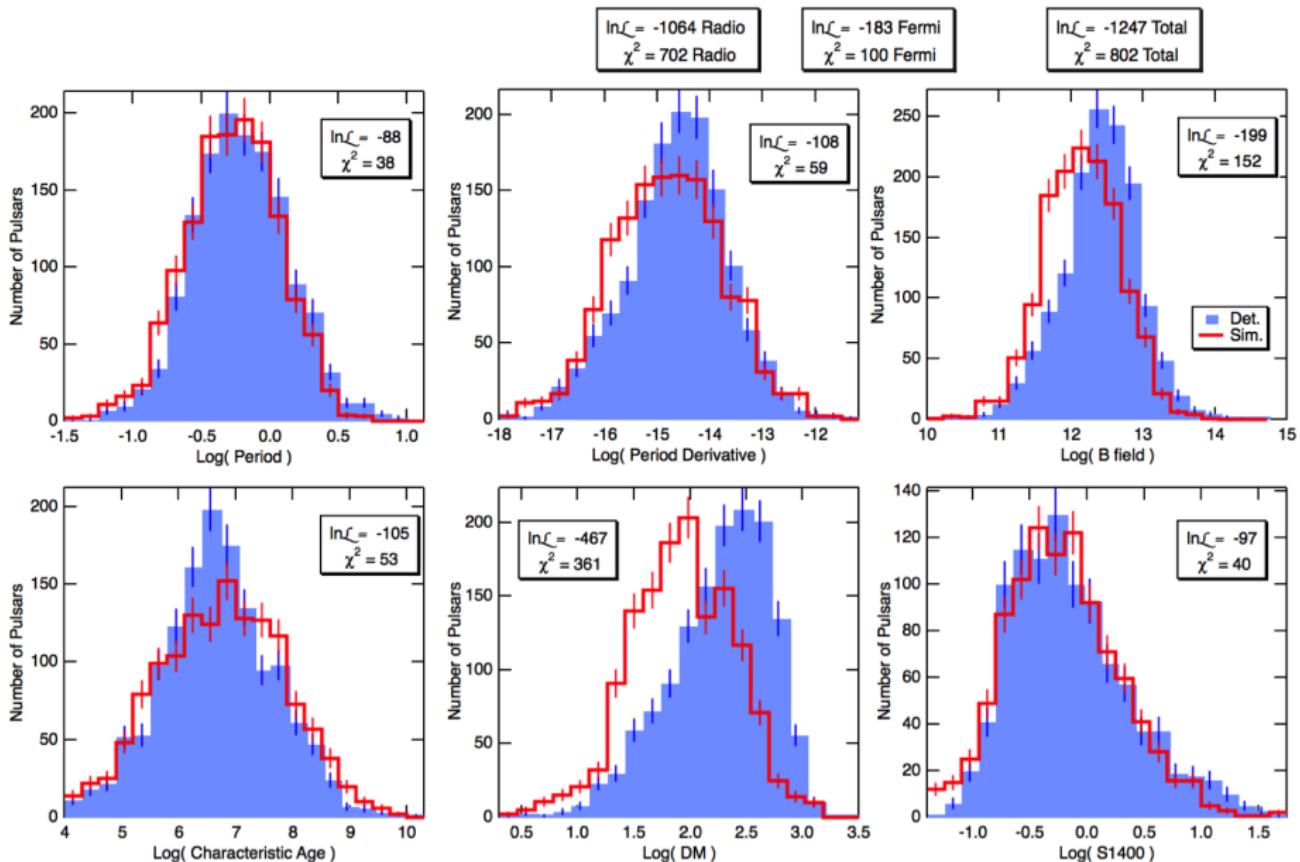
$$\alpha_\gamma = -2.55, \quad \beta_\gamma = 0.63$$

# *B* field power-law decay model

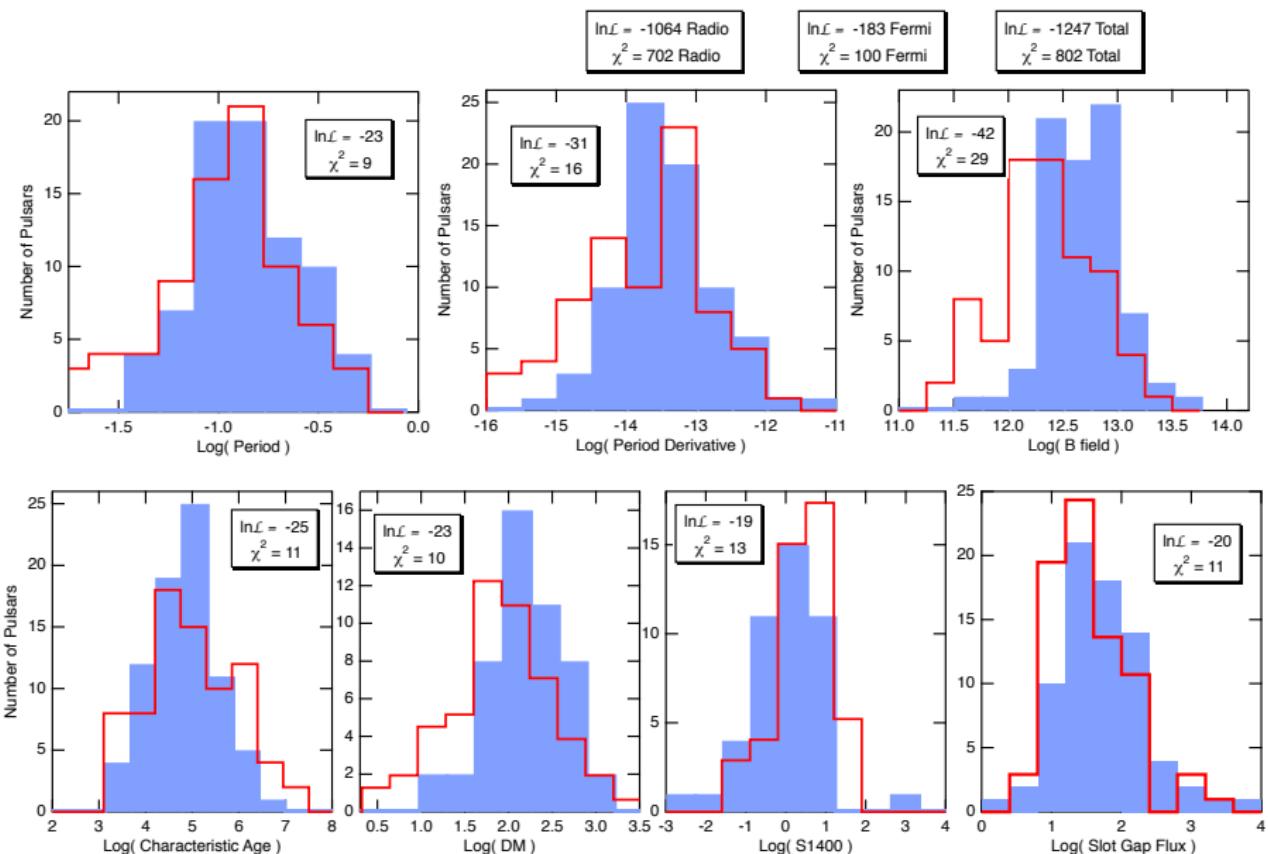


*The Structure and Signals of Neutron Stars, from Birth to Death, March 24 - 28, 2014, Florence, Italy*

# Radio pulsars



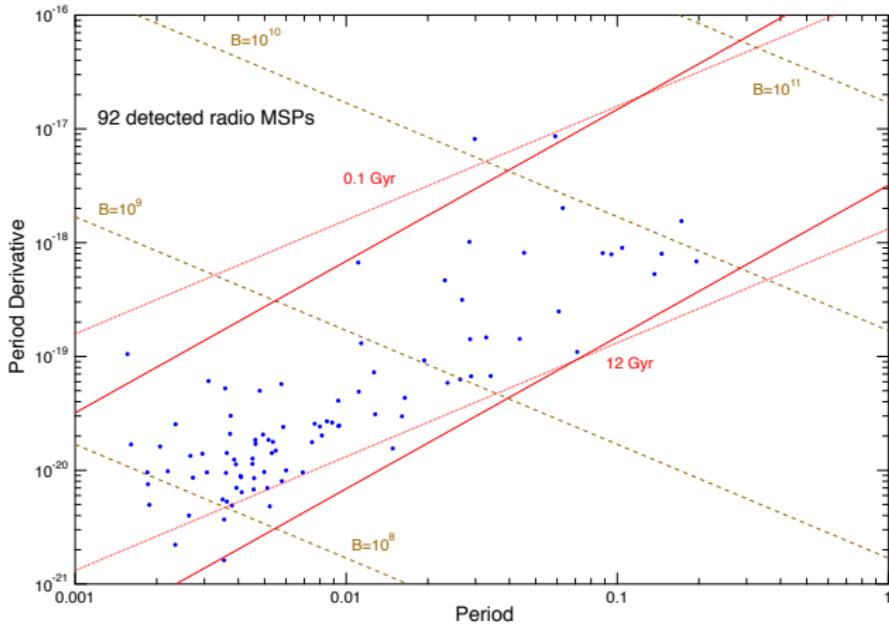
# *Fermi* pulsars



# Population synthesis of millisecond pulsars

## Initial Magnetic Field and Period

- $P(B_8) \propto B_8^{-1.3}$
- $B_{\min} = 0.9B_8$
- $P_{o\min} = 1.3$  ms



Mass accretion lines —  $P_o = 0.18 \times 10^{3\delta/7} \times B_8^{6/7}$  ms where  $\delta$  dithered between 0 and 2.8 — Lamb & Yu (2005)

## MCMC - Radio and $\gamma$ -ray Luminosities

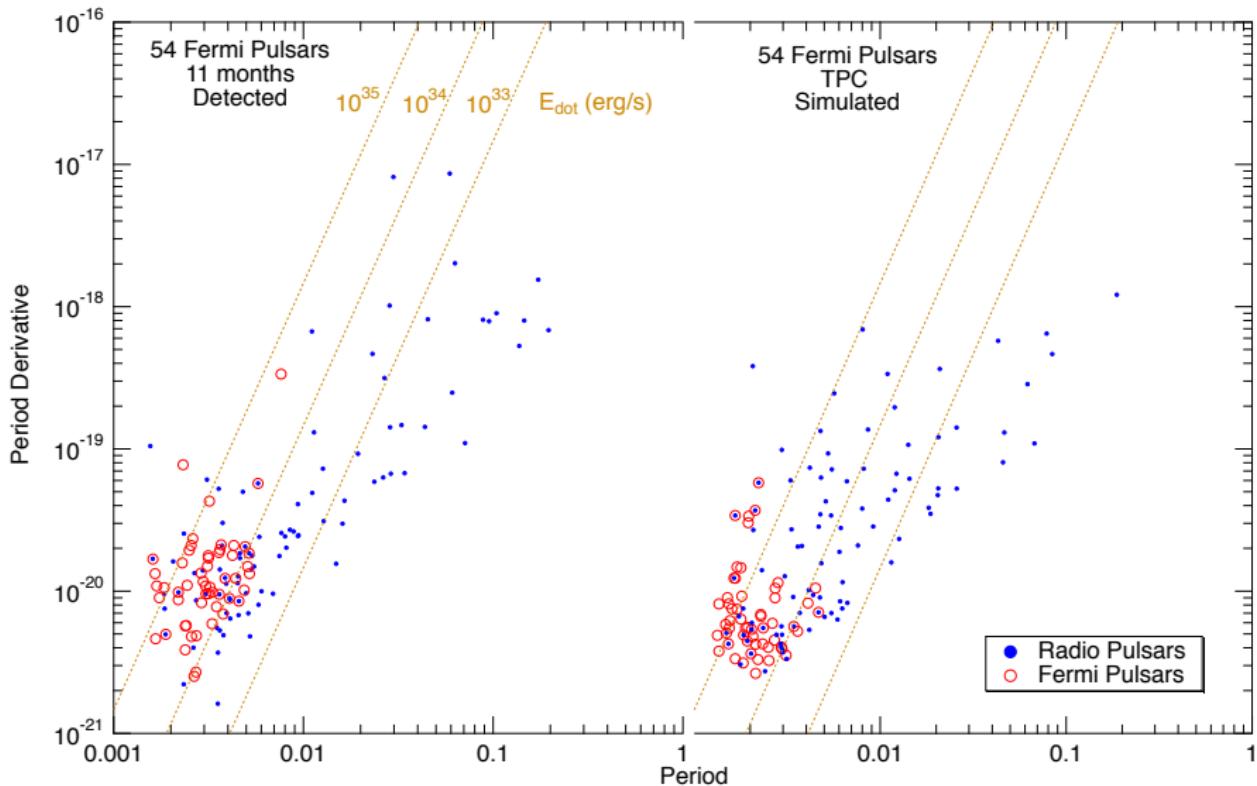
- Assuming empirical radio and  $\gamma$ -ray luminosities as it the case of NPs

$$L \propto P^\alpha \dot{P}^\beta$$

- MCMC searches a 4D model parameter space selecting

$$\alpha_\nu = -1.07 \pm 0.17, \quad \beta_\nu = 0.59 \pm 0.12$$

$$\alpha_\gamma = -2.7 \pm 1.0, \quad \beta_\gamma = 1.1 \pm 0.4$$

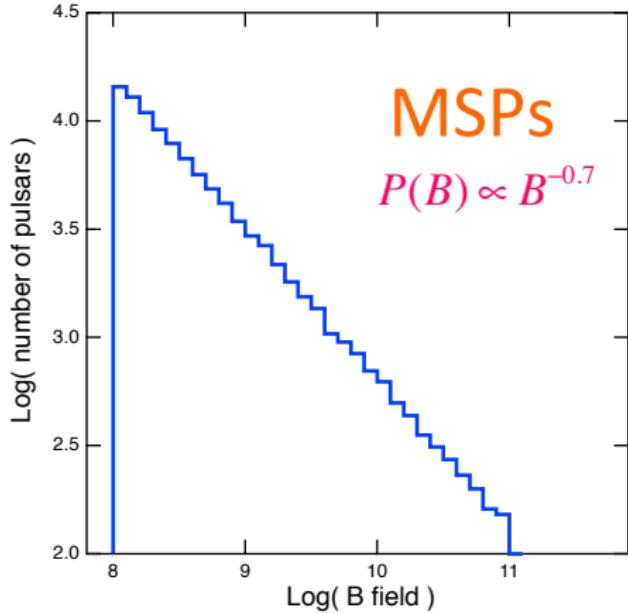
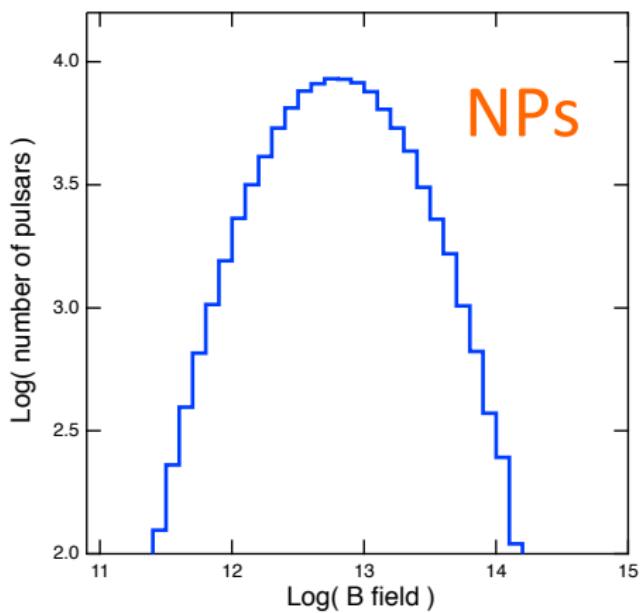


# Summary of Population Synthesis of MSPs

| Catalog | Period    | Detected | Simulated |     |        |      |
|---------|-----------|----------|-----------|-----|--------|------|
|         |           |          | TPC       | OG  | RALTPC | PSPC |
| BSL     | 3 months  | 13       | 29        | 30  | 27     | 29   |
| 1FGL    | 11 months | 54       | 54        | 54  | 54     | 54   |
| 2FGL    | 2 years   | 68       | 76        | 77  | 80     | 80   |
| 3FLG    | 4 years   | 82       | 107       | 106 | 110    | 109  |
|         | 5 years   |          | 119       | 118 | 121    | 122  |
|         | 10 years  |          | 162       | 153 | 170    | 160  |

# The Alicante connection: NPs to MSPs - field decay

How do we go from a log-normal  $B$  distribution (NPs) to a power law  $B$  distribution (MSPs)?



# A natural consequence of the power-law decay of the magnetic field!

$$B = \frac{B_0}{[1 + a \alpha B_0^\alpha t_6]^{1/\alpha}}$$

For large times  $(a \alpha B_0^\alpha t_6) \gg 1$

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We are assuming a constant birth rate, therefore the number of present pulsars at time  $t$  is

$$N = c t_6$$

where  $c$  is the constant birth rate of MSPs —  $c = 4$  to  $5$  per Myr

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$$N = \frac{B^{-\alpha}}{a \alpha c}$$

## Issues based on previous simulations

$$B_{\min}^{\alpha} = \frac{1}{a \alpha t_{6_{\max}}} = (9 \times 10^{-6})^{\alpha}$$

If we assume,  $P(B_8) \propto B_8^{-1.3}$  — as we have in previous simulations

- $\alpha = 1.3$  does not work!

$$a = \frac{1}{1.3 \times (9 \times 10^{-6})^{1.3} \times 12000} = 232 \text{ Myr}^{-1}$$

- So we take  $a = 1 \text{ Myr}^{-1}$  and  $\alpha = 0.7$  for NPs (Viganò et al. 2013) and apply them to MSPs

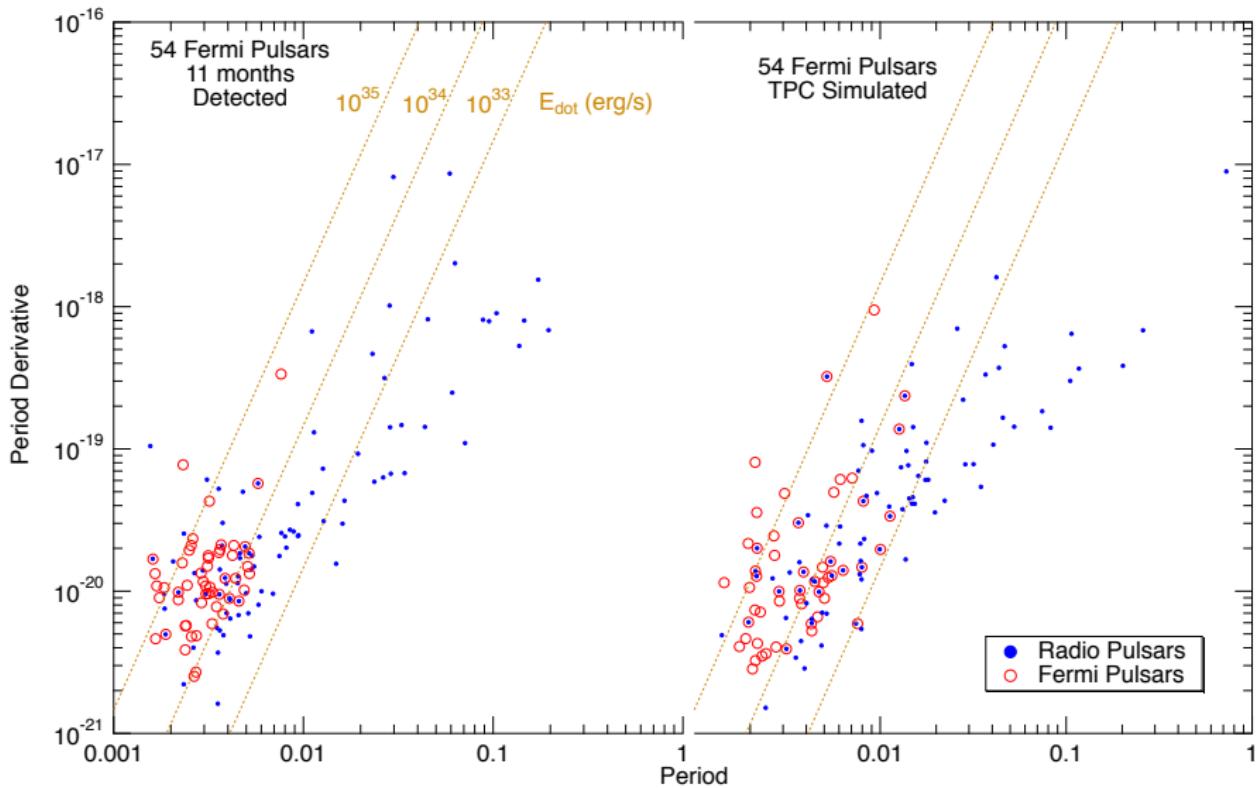
$$t_{6_{\max}} = \frac{1}{0.7 \times 1 \text{ Myr}^{-1} \times (9 \times 10^{-6})^{0.7}} = 4900 \text{ Myr}$$

## Proposed new simulation

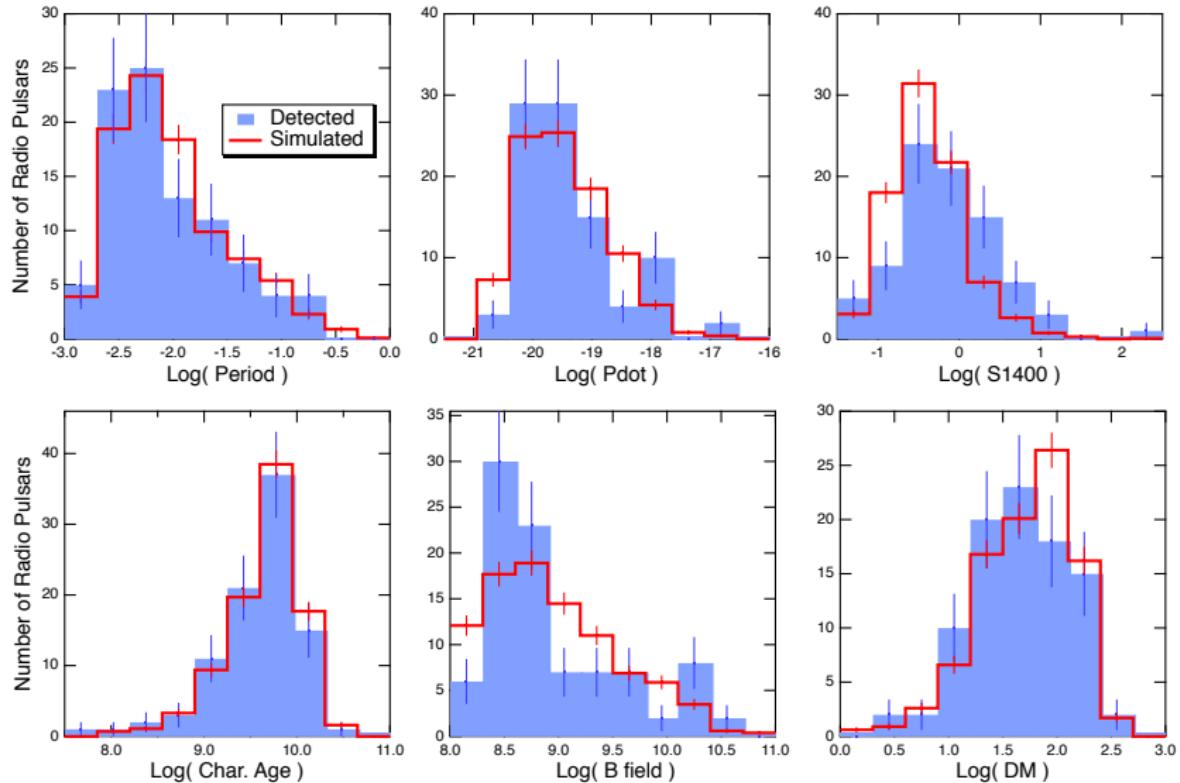
- From NPs  $\alpha = 0.7$  and  $a = 1 \text{ Myr}^{-1}$  to describe Hall induction and Ohmic decay of crustal field
- Switch to  $\alpha = 1.3$  and  $a = 0.01 \text{ Myr}^{-1}$  to describe ambipolar diffusion of the core field — pure guess, but not sensitive
- Theoretical details pending and the code development
- Redo the MSP simulation to test the scenario with

$$P(B_8) \propto B_8^{-0.7}$$

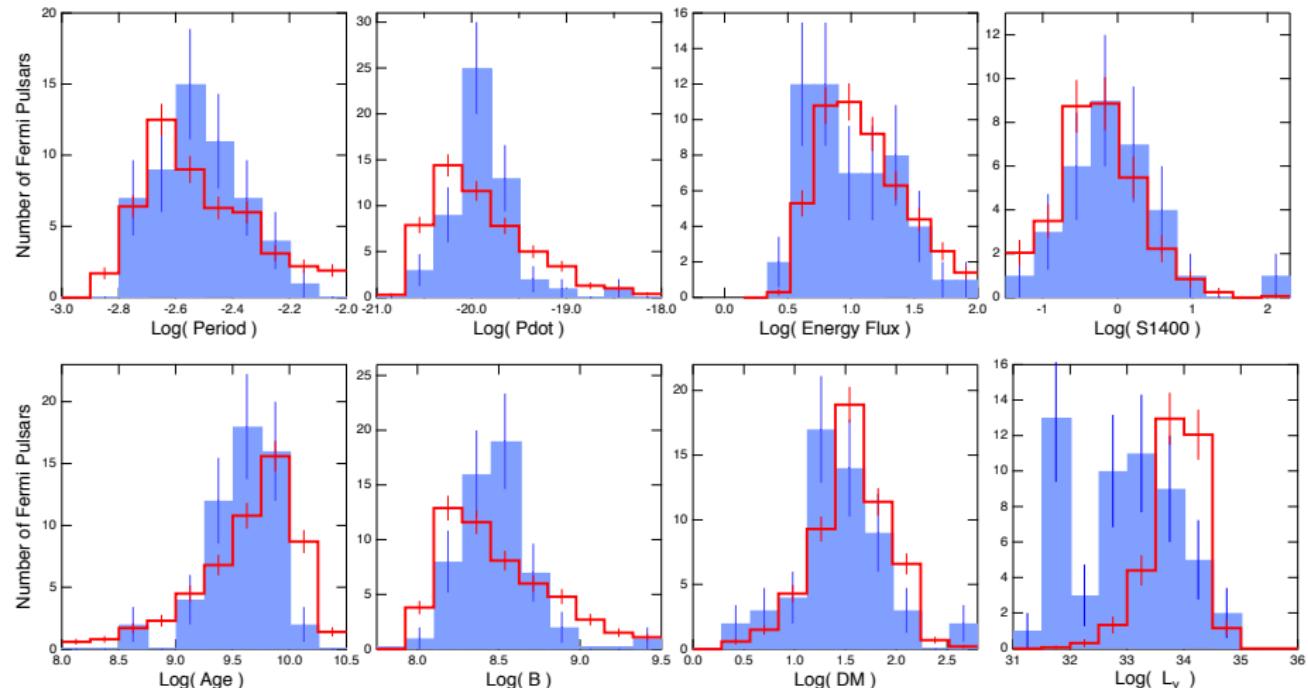
- and  $\alpha = 1.3$  and  $a = 0.01 \text{ Myr}^{-1}$



# Radio pulsars



From  $\alpha_\nu = -1.07$ ,  $\beta_\nu = 0.59$ , to  $\alpha_\nu = -0.94$ ,  $\beta_\nu = 0.05$



From  $\alpha_\gamma = -2.7$ ,  $\beta_\gamma = 1.1$ , to  $\alpha_\gamma = -1.77$ ,  $\beta_\gamma = 0.23$

- Normal pulsars and binary pulsar systems may come from a similar parent distribution of magnetized neutron stars.
- A log-normal  $B$  field distribution of young (normal, isolated) pulsars can be converted into a power-law  $B$  field distribution via standard crustal Hall induction and Ohmic field decay leaving ambipolar diffusion to later play a role within the neutron star core of old dead pulsars that can then be recycled into millisecond pulsars via mass accretion from a stellar companion spinning up the neutron star.
- The cutoff  $B_{\min}$  in the assumed  $B$  field birth distribution of MSPs may be an indication of maximum age and minimum period of active MSPs.
- Detailed theoretical work is pending as well as code development for an adequate full population simulation.